PROCEEDINGS

4th International Congress of Serbian Society of Mechanics

Editors:
Stevan Maksimović
Tomislav Igić
Nataša Trišović

June 4-7, 2013, Vrnjačka Banja, Serbia



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PREFACE

These proceedings contains the papers presented at the Forth (29th Yu) International Congress of Serbian Society of Mechanics held in Vrnjačka Banja during the period 4th – 7th June, 2013.

Theoretical and Applied Mechanics is a subject of great importance in the developing of science and technology. The aim of the Congress is to provide a forum to exhibit the progress in this field during the past years and a place to further the interaction between of modern theoretical and applied mechanics, as well as modern engineering sciences.

The papers, contributed by authors from all around the globe, have been separated into 7 sections which cover the main areas of interest: 'Plenary Lectures', Section A, Section B, Section C, Section D and two Mini-symposia.

We would like to express our gratitude to all members of the Scientific Committee and also to the participants for their engagement in organizing of the Congress, including the preparation of manuscripts to be published in the Journal Theoretical and Applied Mechanics, Scientific Technical Review and Journal of Serbian Society for Computational Mechanics.

It gives us great pleasure to express our deep appreciation for the great long-standing support that Prof. Dr. Nikola Hajdin, President of the Serbian Academy of Sciences and Arts, has given to the promotion of all aspects of theoretical and applied mechanics in Serbia.

Last, the Congress organizing committee wishes to acknowledge the collaboration of the Ministry of Education, Science and Technological Development of the Republic of Serbia, Serbian Academy of Sciences and Arts, Municipality Vrnjačka Banja and many supporting members of the Serbian Society of Mechanics.

S. MAKSIMOVIĆ & T. IGIĆ June, 2013

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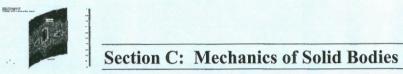


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STRESS INTEGRATION OF MATSUOKA-NAKAI CONSTITUTIVE MODEL USING INCREMENTAL PLASTICITY METHOD

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Abstract. This paper presents the integration of constitutive relations for Matsuoka-Nakai non-associative constitutive model using incremental plasticity method. The basic idea of this method is the calculation of plastic matrix thus, the calculation of the stress increment needs only total strain increment. Presented constitutive model consists of a yield surface similar to Mohr-Coulomb yield surface, but completely smooth, avoiding corners which can be problematic in numerical calculations. The yield surface divides stress 3D space into the pure elastic region and pure plastic region. In addition to the basic equations for stress integration, the paper presents the development of elastic-plastic constitutive matrix and algorithm for the implementation in the program for finite elements. Presented procedure is implemented in the PAK software. Verification is performed using examples from the literature and compared with the results obtained using Mohr-Coulomb constitutive model.

1. Introduction

Stress integration represents calculation of stress change during an incremental step, corresponding to strain increments in the step. It is in essence the incremental integration of inelastic constitutive relations to trace the history of material deformation. The stress integration is an important ingredient in the overall finite element inelastic analysis of structures. It is important that the integration algorithm accurately reproduces the material behavior since the mechanical response of the entire structure is directly dependent on this accuracy. The algorithm should be also computationally efficient because the stress integration is performed at all integration points. For general applications, this computational procedure should be robust, providing reliable results under all possible loading conditions. In this paper we present a formulation of the computational algorithm for the Matsuoka-Nakai (MN) constitutive model [1] using an incremental plasticity approach (IPM). The solutions were compared with the results obtained using Mohr-Coulomb (MC) model [2], as well as with the results obtained using concrete constitutive model.

In the next section we present formulation the MN constitutive model, followed by the derivation of the elastic-plastic constitutive matrix in general associated plasticity. Then, the general relations are implemented in MN constitutive model.

2. Matsuoka-Nakai constitutive model formulation

Matsuoka-Nakai constitutive model is based on experimental results of soil material research. Failure surface of this constitutive model is described using stress invariants in the following form:

$$f = I_3 + \frac{\cos^2 \phi}{9 - \sin^2 \phi} I_1 I_2 = 0 \tag{1}$$

where I_1 , I_2 and I_3 are first, second and third stress invariants, respectively [3]. Parameter ϕ represents the angle of material internal friction. Failure surface of this constitutive model in deviatoric plane is presented in Figure 1.

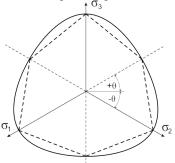


Figure 1 Mohr-Coulomb and Matsuoka-Nakai failure surface in deviatoric plane

As seen in Figure 1, for axisymmetric stress state, failure surface of Matsuoka-Nakai model matches the failure surface of Mohr-Coulomb model. However, unlike the failure surface of Mohr-Coulomb model, failure surface of Matsuoka-Nakai model is completely smooth, which is more suitable from the aspect of numerical solving.

In the case of non-associative yield condition, failure surface and the plastic potential surface of this model are different. In that case, plastic potential surface is presented by the following equation:

$$g = I_3 + \frac{\cos^2 \psi}{9 - \sin^2 \psi} I_1 I_2 = 0 \tag{2}$$

where ψ represents dilatation angle of material.

In the case of cohesive materials, failure surface (1) and plastic potential surface (2) are formulated using modified stress tensor [4] as:

$$\sigma_{ii}^* = \sigma_{ii} - \sigma_0 \delta_{ii} \tag{3}$$

where σ_{ij} represents effective stress tensor, δ_{ij} is Kronecker delta symbol, whereas σ_0 is defined as:

$$\sigma_0 = c \ ctg \ \phi \tag{4}$$

where c is material cohesion and ϕ is previously mentioned internal fraction angle.

3. Elastic-plastic constitutive matrix using IPM

Elastic-plastic constitutive models are described through elastic-plastic constitutive relations. In incremental plasticity theory, stress is directly proportional to strain up to reaching yield stress. After reaching yield stress, strain increment can be divided into elastic and plastic part [5]:

$$\left\{de\right\} = \left\{de^{E}\right\} + \left\{de^{P}\right\} \tag{5}$$

Only elastic part of strain causes the stress change thus the stress increment can be formulated as:

$$\{d\sigma\} = \left\lceil C^E \right\rceil \left\{ de^E \right\} \tag{6}$$

where $\lceil C^E \rceil$ is elastic constitutive matrix. Substituting (5) in (6) the following is obtained:

$$\{d\sigma\} = \left[C^{E}\right] \left(\{de\} - \left\{de^{P}\right\}\right) \tag{7}$$

In the case of using elastic-plastic constitutive models, yield function is the stress state function, therefore the increment of its change can be formulated as:

$$df = \left\{\frac{\partial f}{\partial \sigma}\right\}^T \left\{d\sigma\right\} = 0 \tag{8}$$

In incremental plasticity theory it is necessary that the failure function is in every time step less or equal to zero (neutral loading condition).

Implicit stress integration implies the increment of plastic strain in the normal direction on the plastic potential surface, which can be formulated as:

$$\left\{ de^{P} \right\} = d\lambda \left\{ \frac{\partial g}{\partial \sigma} \right\} \tag{9}$$

where $d\lambda$ is positive scalar which is to be calculated and plastic potential function g is the stress state function. Substituting the plastic strain increment (9) in(7) and using (8), it is obtained:

$$df = \left\{ \frac{\partial f}{\partial \sigma} \right\}^{T} \left(\left[C^{E} \right] \left\{ de \right\} - d\lambda \left[C^{E} \right] \left\{ \frac{\partial g}{\partial \sigma} \right\} \right) = 0$$
 (10)

Plastic parameter $d\lambda$ can be calculated from equation (10) as:

$$d\lambda = \frac{\left\{\frac{\partial f}{\partial \sigma}\right\}^T \left[C^E\right] \left\{de\right\}}{\left\{\frac{\partial f}{\partial \sigma}\right\}^T \left[C^E\right] \left\{\frac{\partial g}{\partial \sigma}\right\}}$$
(11)

Finally, using parameter $d\lambda$ from (11) stress increment $\{d\sigma\}$ is obtained using (7) and (9) in the function of total strain increment:

$$\{d\sigma\} = \lceil C^{EP} \rceil \{de\} \tag{12}$$

where term $\lceil C^{\textit{EP}} \rceil$ represents elastic-plastic constitutive matrix.

4. Stress integration of Matsuoka-Nakai model

The derivative of the MN failure function (1) with respect to stresses can be calculated using the chain rule:

$$\left\{\frac{\partial f}{\partial \sigma}\right\}^{T} = \frac{\partial f}{\partial I_{1}} \left\{\frac{\partial I_{1}}{\partial \sigma}\right\}^{T} + \frac{\partial f}{\partial I_{2}} \left\{\frac{\partial I_{2}}{\partial \sigma}\right\}^{T} + \frac{\partial f}{\partial I_{3}} \left\{\frac{\partial I_{3}}{\partial \sigma}\right\}^{T}$$

$$(13)$$

The derivative of plastic potential function (2) can be calculated likewise:

$$\left\{ \frac{\partial g}{\partial \sigma} \right\}^{T} = \frac{\partial g}{\partial I_{1}} \left\{ \frac{\partial I_{1}}{\partial \sigma} \right\}^{T} + \frac{\partial g}{\partial I_{2}} \left\{ \frac{\partial I_{2}}{\partial \sigma} \right\}^{T} + \frac{\partial g}{\partial I_{3}} \left\{ \frac{\partial I_{3}}{\partial \sigma} \right\}^{T} \tag{14}$$

Certain derivatives of the failure function from (13) are:

$$\frac{\partial f}{\partial I_1} = \frac{\cos^2 \phi}{9 - \sin^2 \phi} I_2, \quad \frac{\partial f}{\partial I_2} = \frac{\cos^2 \phi}{9 - \sin^2 \phi} I_1, \quad \frac{\partial f}{\partial I_3} = 1$$
 (15)

whereas derivatives of plastic potential function from (14) are:

$$\frac{\partial g}{\partial I_1} = \frac{\cos^2 \psi}{9 - \sin^2 \psi} I_2, \quad \frac{\partial g}{\partial I_2} = \frac{\cos^2 \psi}{9 - \sin^2 \psi} I_1, \quad \frac{\partial g}{\partial I_3} = 1$$
 (16)

Using equations (13) to (16) elastic-plastic constitutive matrix can be calculated and then the stress increment (12) as well. According to this theory, algorithm for stress integration can be formed using incremental plasticity theory as given in Table 1.

Table 1 Stress integration algorithm using IPM

A. Known:
$$\{{}^{t+\Delta t}e\}$$
, $\{{}^{t}e\}$, $\{{}^{t}\sigma\}$, $\{{}^{t}e^{p}\}$
 $\{d\sigma\} = [C^{E}]\{de^{E}\} = [C^{E}](\{{}^{t+\Delta t}e\} - \{{}^{t}e\})$, $\{{}^{t+\Delta t}\sigma\} = \{{}^{t}\sigma\} + \{d\sigma\}$

B. Check the solution IF (F < 0) elastic solution (GOTO E)

 $IF (F \ge 0)$ elastic – plastic solution (CONTINUE)

$$\left\{\frac{\partial f}{\partial \sigma}\right\}, \quad \left\{\frac{\partial g}{\partial \sigma}\right\} \quad \text{using} \quad (13) \text{ and } (14) \text{ and } d\lambda \text{ using } (11)$$

C. Local iterations

 $d\lambda$ correction

$$\text{calculation } \left\{ de^P \right\} = d\lambda \left\{ \frac{\partial g}{\partial \sigma} \right\} \; , \; \left\{ de^E \right\} = \left\{ de \right\} - \left\{ de^P \right\} \; , \; \left\{ d\sigma \right\} = \left[C^E \right] \left\{ de^E \right\}$$

calculation $f(\sigma_{ij})$ using (1)

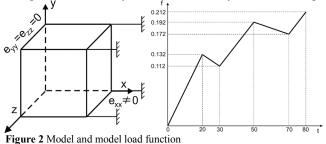
D. Check: $IF(ABS(f) \ge toll)$ return to C

$$\left\{{}^{t+\Delta t}e^{P}\right\} = \left\{{}^{t}e^{P}\right\} + \left\{de^{P}\right\}$$

E. End: $\{{}^{t+\Delta t}\sigma\},\{{}^{t+\Delta t}e^p\}$

5. Verification examples

Verification of developed algorithm of MN constitutive model is performed on the example of one-axis model compression. Model consists of single element loaded towards (Figure 2a). Model was loaded using prescribed displacement on one element face, using the function in Figure 2b. Boundary conditions of analyzed model are presented in Figure 2a.



The results are compared with the results obtained in analysis of the same example using Mohr-Coulomb constitutive model (Figure 3).

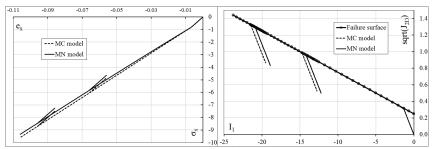


Figure 3 Axial stress as a function of strain and second deviatoric stress invariant as the function of first stress invariant

Analyzing results in Figure 3 can be noticed a relatively small deviation of results obtained using MN model compared with the results obtained using MC model. Also, the MN model is more stable in numerical terms, because the yield surface smooth, unlike the MC model.

Second example represents using of MN constitutive model in the analysis of reinforced concrete beam failure [6]. Model geometry, boundary conditions and load are presented in Figure 4. Eight-node brick finite elements for concrete beam modeling are used, whereas beam elements of appropriate dimensions are used for reinforcement. Model was loaded using prescribed displacement.

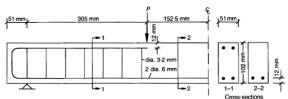


Figure 4 Load scheme, boundary condition and geometry of concrete beam

Comparison of results obtained using MN constitutive model with the results obtained using concrete constitutive model are presented in Figure 5.

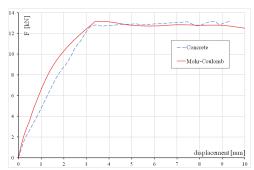


Figure 5 Force-displacement result

Analyzing results in the Figure 5, we can notice good matching of the solutions obtained using MN model with the solutions obtained using concrete constitutive model. It can be concluded that the MN model can be used to simulate the mechanical behavior of concrete structures. Mechanical characteristics of MN model for simulating the concrete behavior can be obtained by fitting the failure surfaces of these two models.

6. Conclusion

The results of the presented constitutive model are compared with the results obtained by Mohr-Coulomb constitutive model and as it can be seen, these two models provide very similar results. Results of MN model are also compared with the results obtained using concrete constitutive model. Analyzing the results a good matching is noticed. The advantage of the presented computational procedure is a general formulation, which can be applied to various yield functions, with yield function expressed in terms of stress invariants. Also, this procedure can be implemented in an explicit integration scheme (no yield condition check D in Table 1).

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